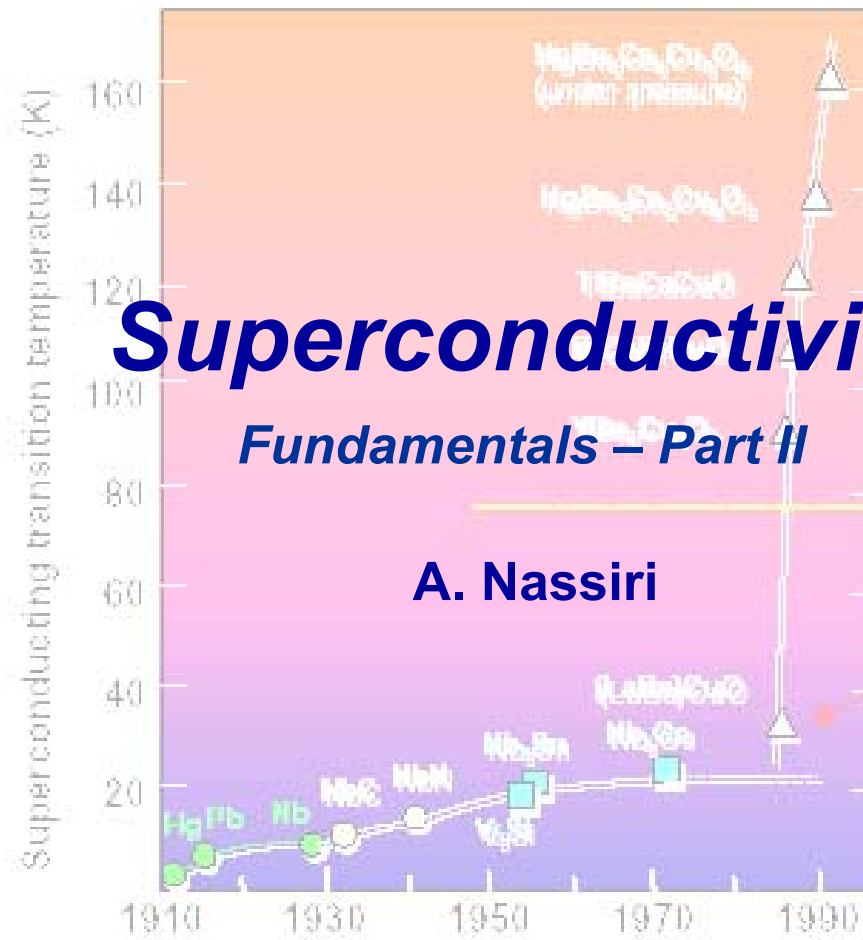
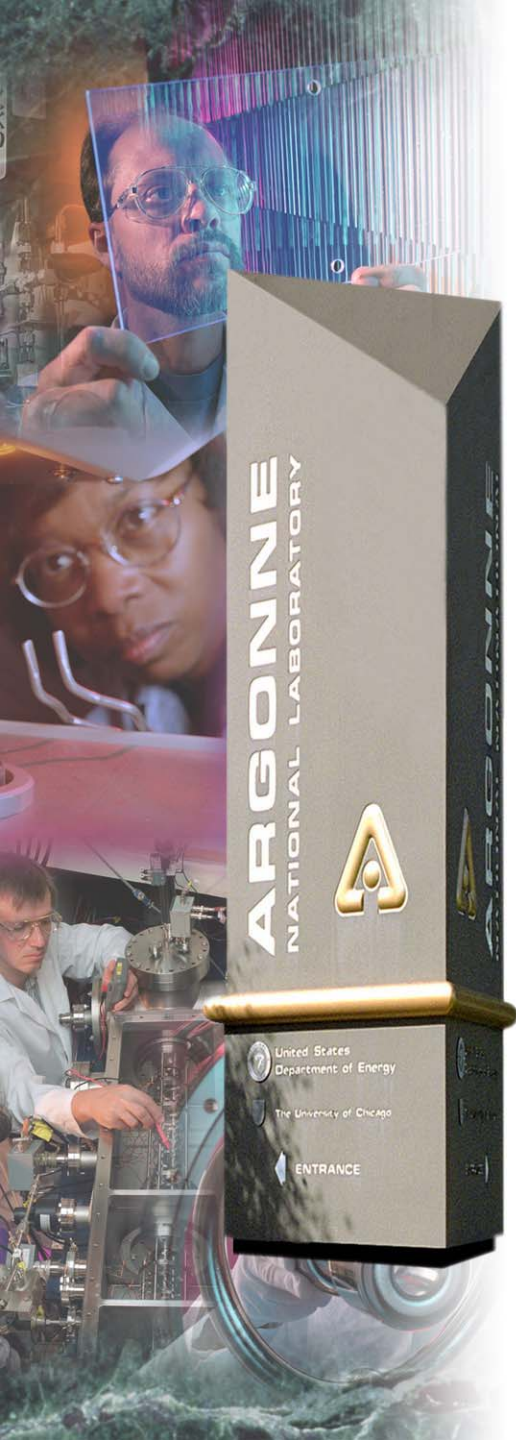




"I WAS HOPING TO SEE MORE EVIDENCE
OF THE FLUX-PINNING SITES."



Superconductivity

Fundamentals – Part II

A. Nassiri

Argonne National Laboratory

Thursday, February 17, 2005



A U.S. Department of Energy
Office of Science Laboratory
Operated by The University of Chicago



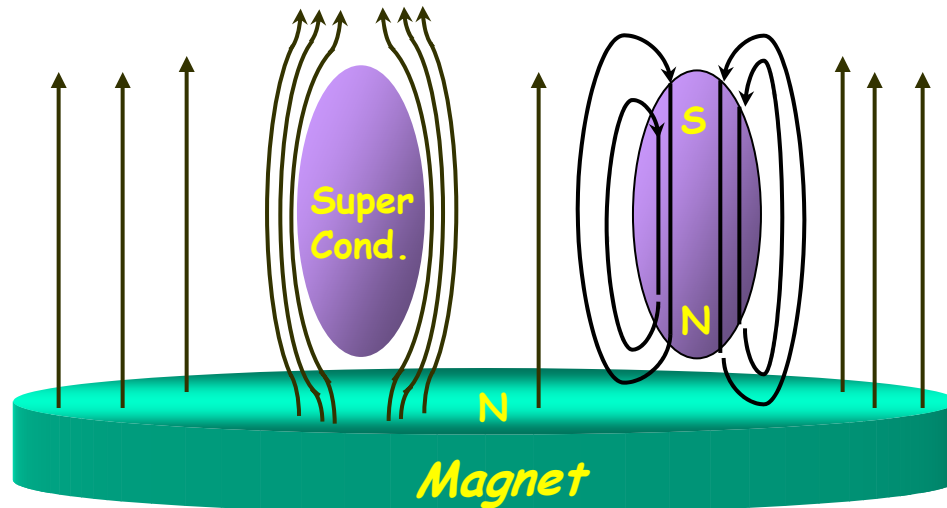
OUTLINE

- **Meissner Effect**
- **Type I superconductor**
- **Type II superconductor**
- **Resistivity**
- **Accelerating gradient limitations**
- **APS SR RF Options**



Properties of Superconductors

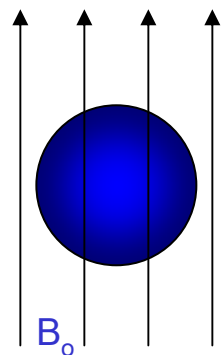
- Zero electrical resistivity ($\rho=0$) -> persistent current
- Meissner effect
- Critical field, critical current
- Penetration depth, coherent length



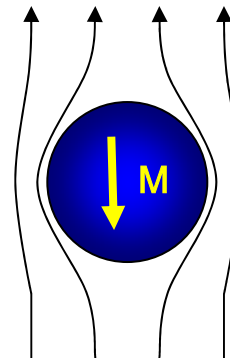
Properties of Superconductors

Meissner effect (Meissner & Ochsenfeld – 1933)

Magnetic field is completely expelled by a superconductor (perfect diamagnetism)



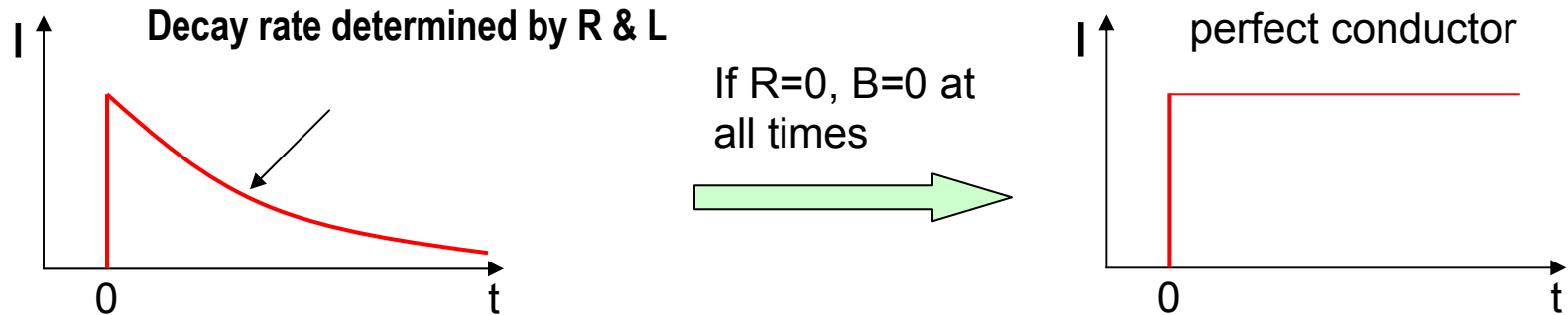
$T > T_c$



$T < T_c$

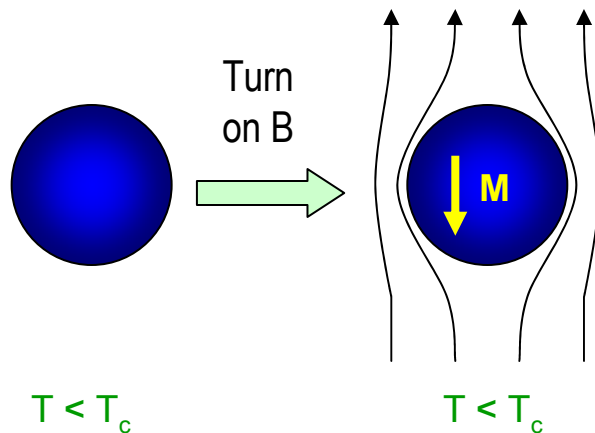
Properties of Superconductors

- Magnetic field decays. Can be viewed as an RL circuit.

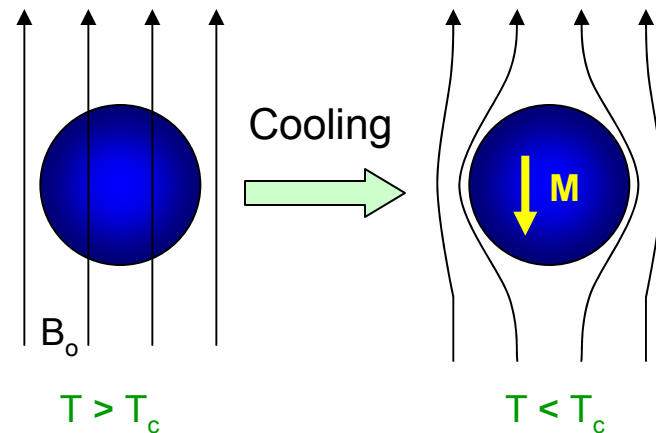


- There is more to superconductor than zero resistance!

Zero field cooling



Field cooling



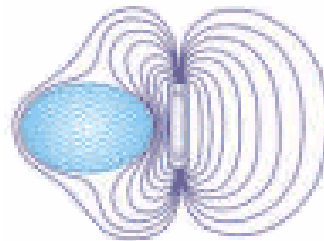
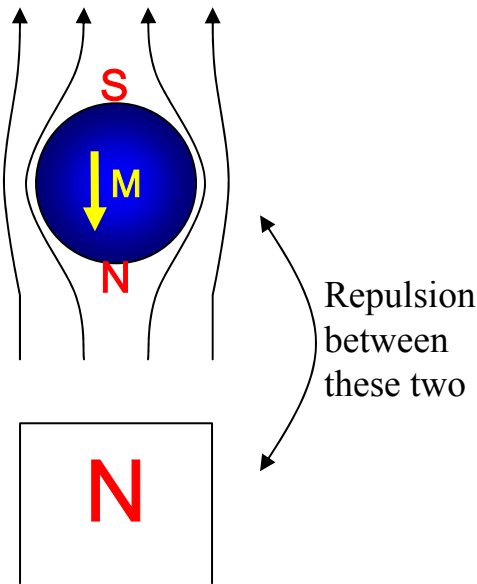
Properties of Superconductors

● Note that just having zero resistance can not explain the field cooling case since there is no change in the external magnetic field upon cooling. Yet magnetic field is still expelled.

● **Meissner effect** is a quantum mechanical effect

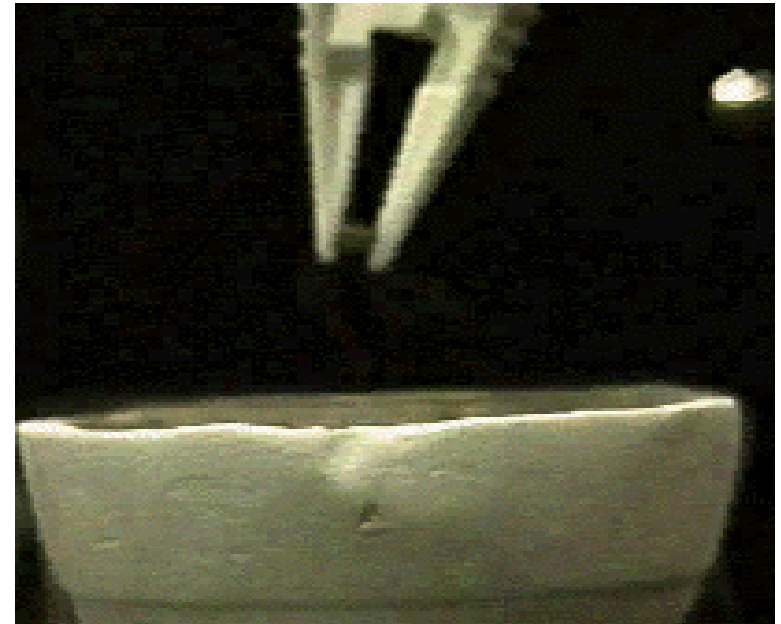
● $\mathbf{B} = 0$ inside the superconductor and ($\mathbf{j}=0$) according to Maxwell's equations. Therefore, there is no current in the superconductor. \Rightarrow Current flows only on the surface.

● Perfect diamagnetism gives a repulsion force \Rightarrow Levitation



Flux exclusion

Perfect diamagnetism prevents magnetic field from penetrating a pure superconductor



induced currents

SC periodic table

KNOWN SUPERCONDUCTIVE ELEMENTS																		
1A																0		
1	H															2	He	
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
3	11 Na	12 Mg	III B	IV B	V B	VI B	VII B	VII		IB	II B	13 Al	14 Si	15 P	16 S	17 Cl	18 Ar	
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
6	55 Cs	56 Ba	57 *La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
7	87 Fr	88 Ra	89 +Ac	104 Rf	105 Ha	106	107	108	109	110	111	112						

■ BLUE = AT AMBIENT PRESSURE

■ GREEN = ONLY UNDER HIGH PRESSURE

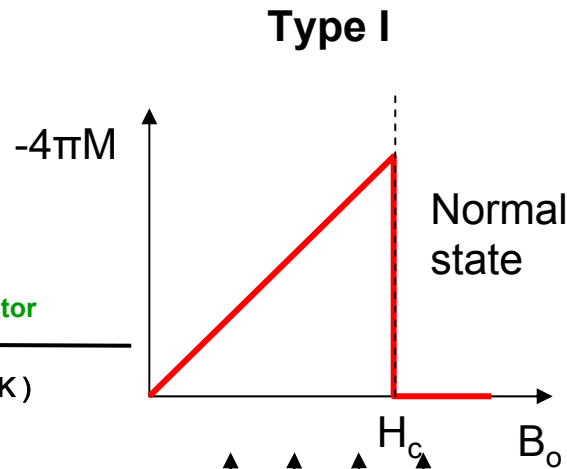
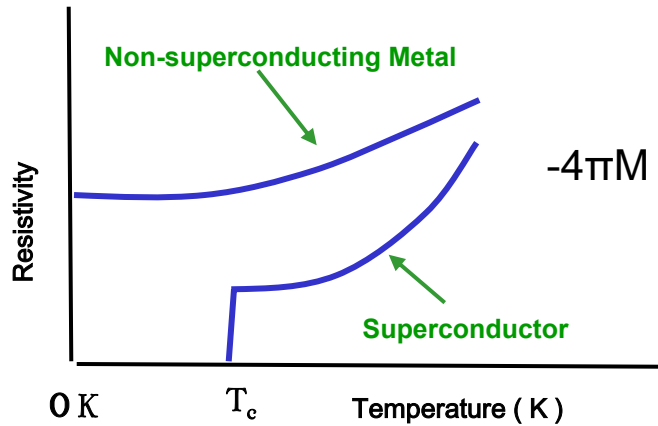
SUPERCONDUCTORS.ORG

* Lanthanide Series

+ Actinide Series

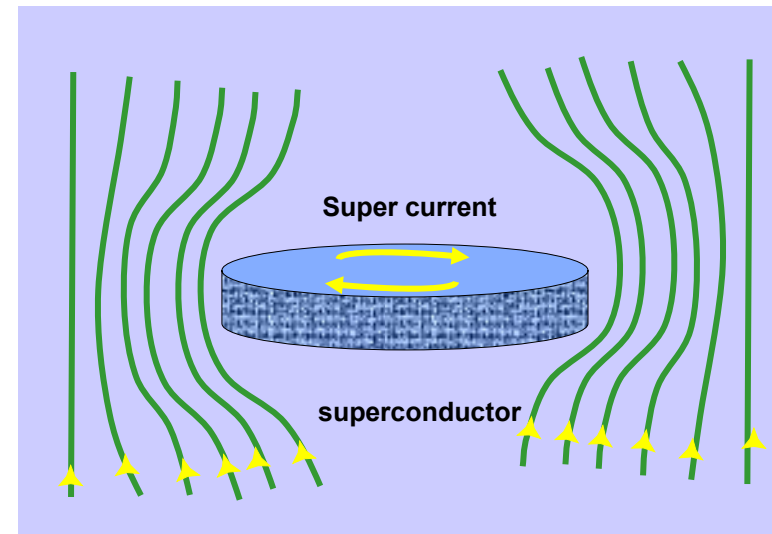
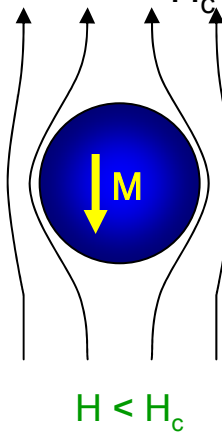
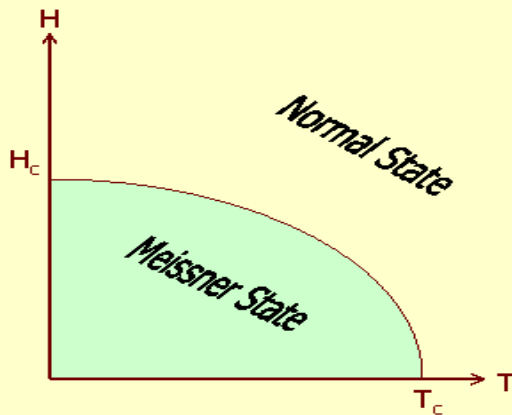
58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr

Type I superconductor



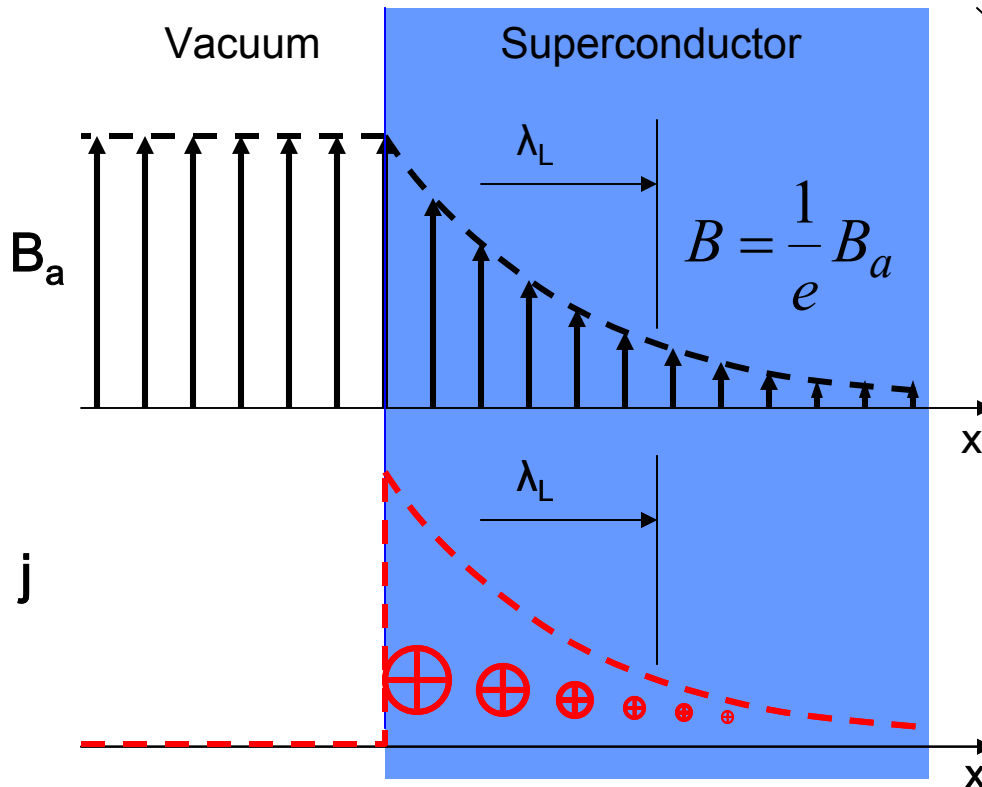
Lead (Pb)	7.196 K
Mercury (Hg)	4.47 K
Tin (Sn)	4.15 K
Aluminum (Al)	1.38 K
Gallium (Ga)	1.175 K
Molybdenum (Mo)	1.083 K
Titanium (Ti)	0.49 K
Uranium (U)	0.40 K
Hafnium (Hf)	0.20 K
Iridium (Ir)	0.128 K
Beryllium (Be)	0.1125 K
Tungsten (W)	0.023 K (SRM 768)

Theory: BCS & Ginsburg-Landau Phase Diagram

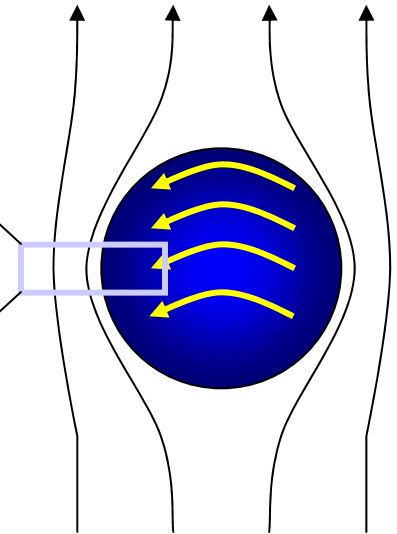


Properties of Superconductors

4. Penetration depth, coherence length



$$B_{inside} = B_a e^{-x/\lambda_L}$$



Note that $\lambda_L^2 = \frac{mc^2}{4\pi e^2 n_s}$

where n_s is the superconducting electron density. If n_s increases, λ_L becomes shorter. In general, pure materials have short λ_L because their n_s is high.



Type II superconductor

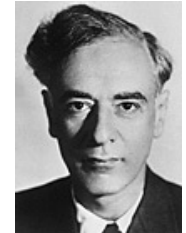
Theory: Abrikosov & Ginsburg-Landau



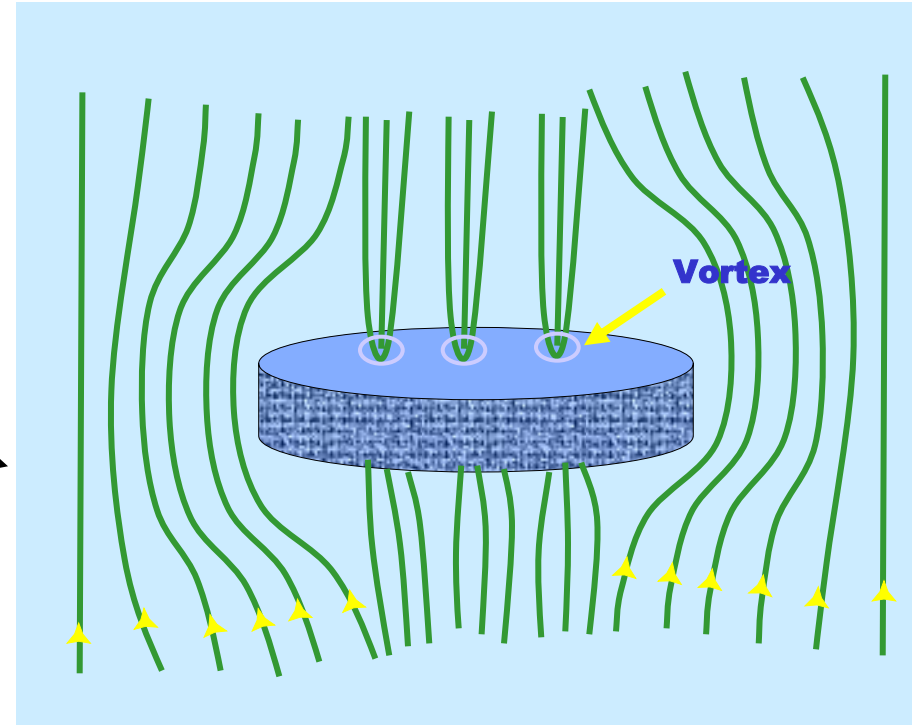
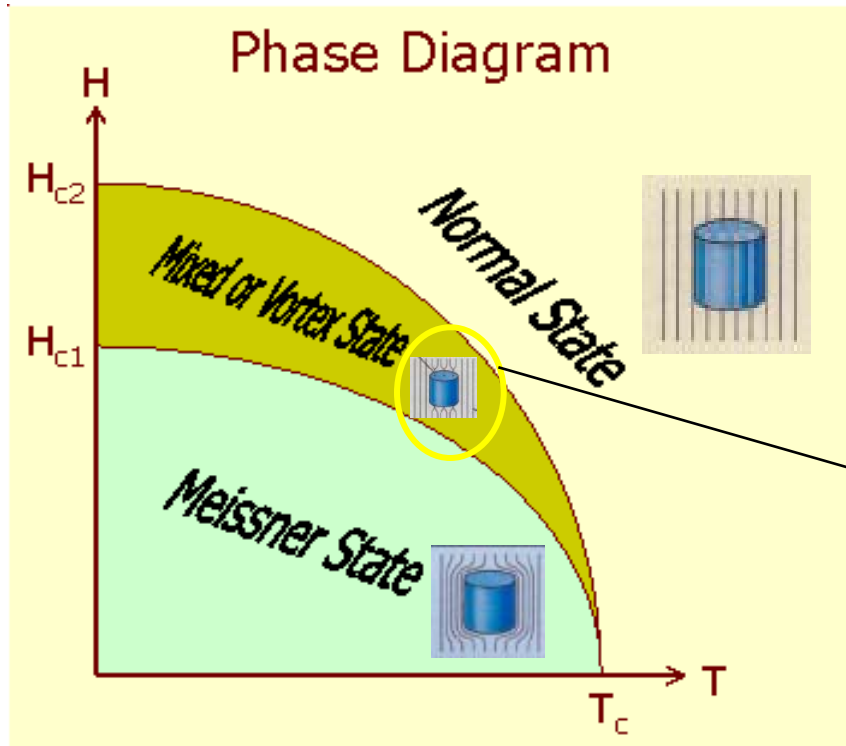
2003



2003



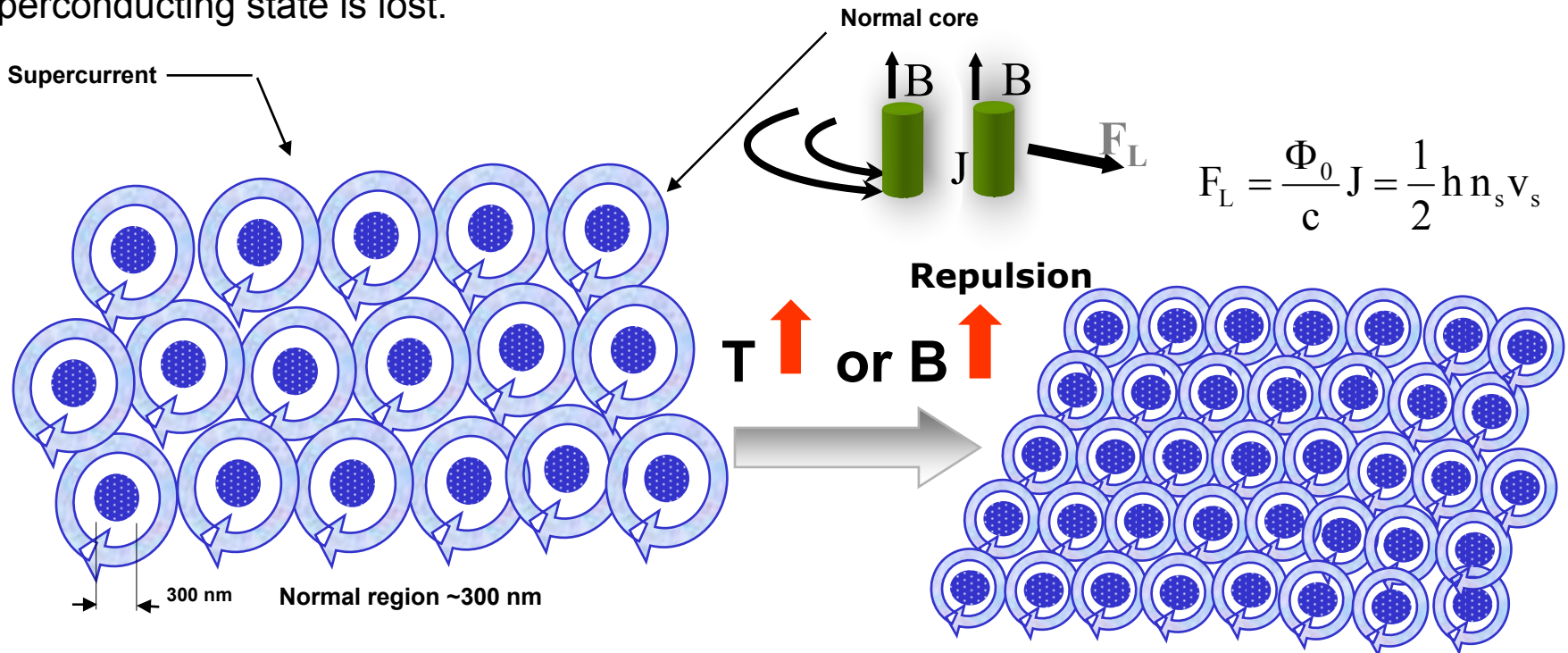
1962



Type II superconductors

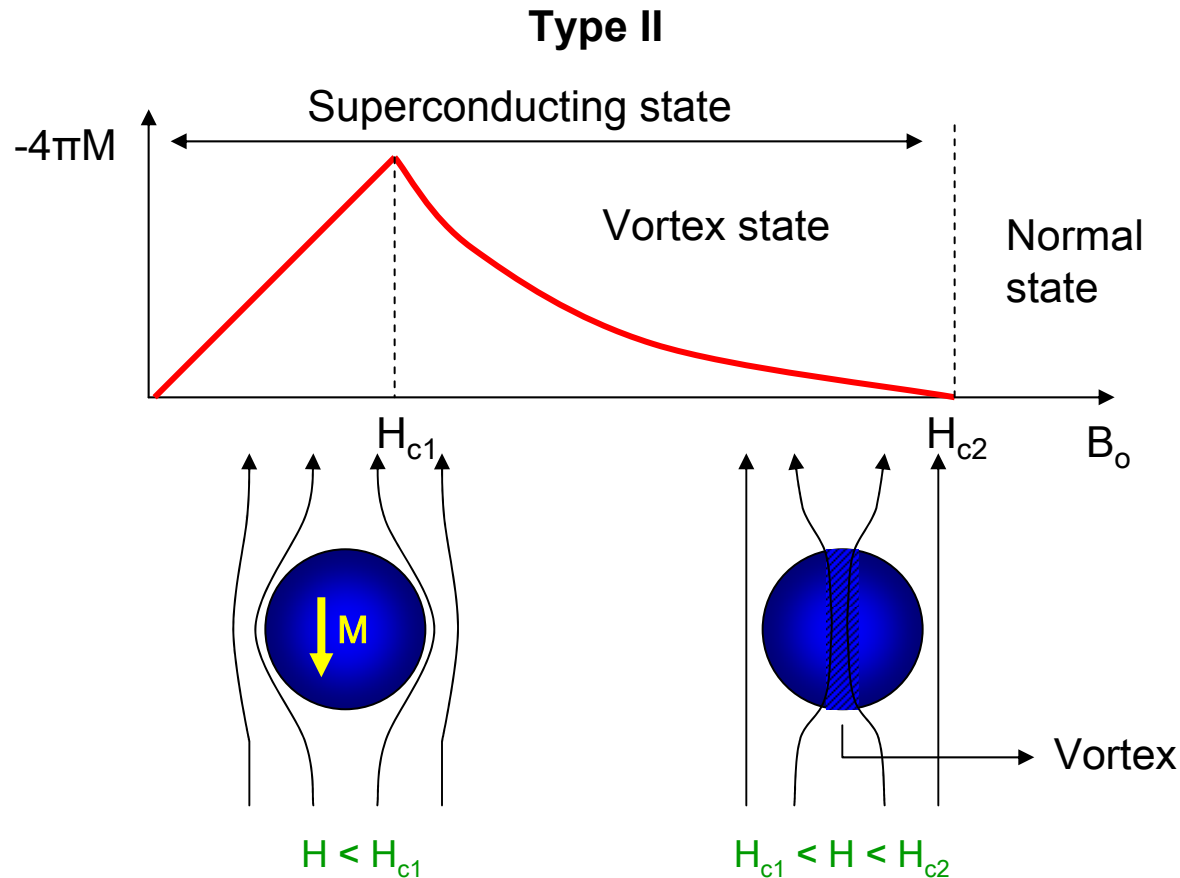
Vortex State for Superconductors

Type II superconductors usually exist in a vortex state with normal cores surrounded by superconducting regions. This allows magnetic field penetration. As their critical temperatures are approached, the normal cores are more closely packed and eventually overlap as the superconducting state is lost.



Properties of Superconductors

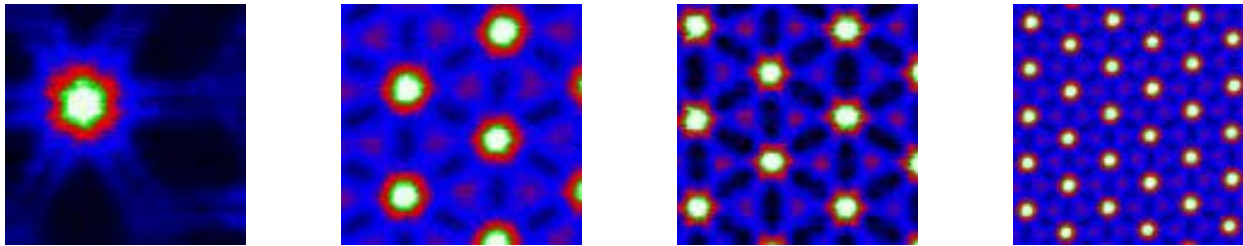
3. Critical field (H_c), critical current density (J_c)



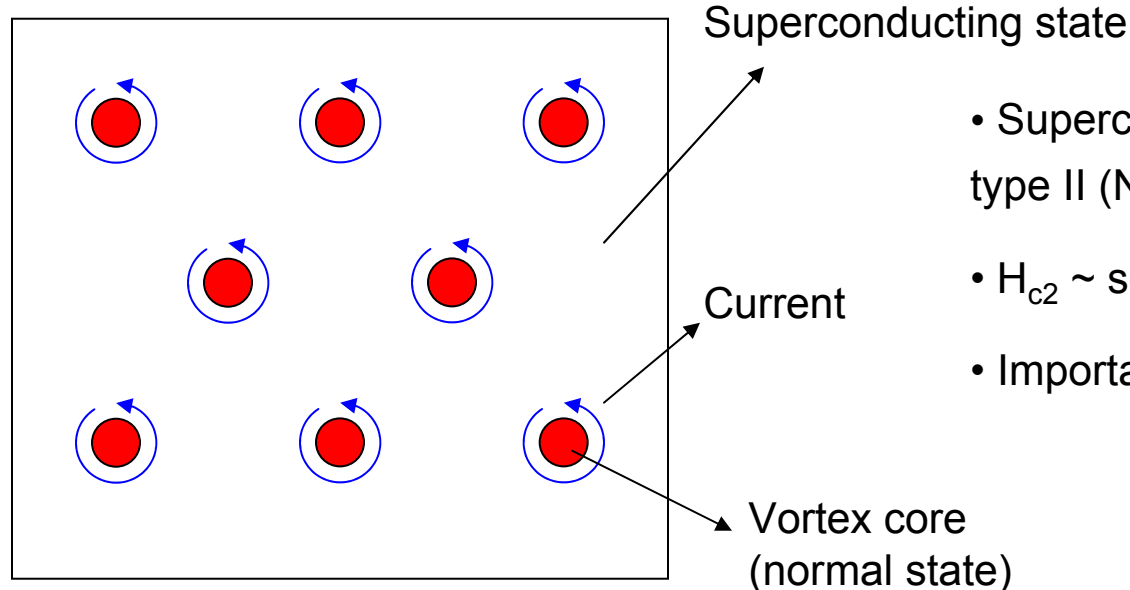
Properties of Superconductors

3. Critical field (H_c), critical current density (J_c)

● Vortex states (STM image on NbSe₂)



B = 500 Gauss \longrightarrow 2 T



- Superconductors with high T_c 's are type II (Nb, Nb₃Se, Nb₃Ge, ...)
- $H_{c2} \sim$ several Tesla
- Important for application



NC vs. SC

Normal Conducting Metal : Skin Depth

$$\delta_s = \sqrt{\frac{2}{\omega * \mu * \sigma}} = \sqrt{\frac{\rho}{\pi * f * \mu}}$$

where:

- μ = permeability ($4\pi * 10^{-7}$ H/m), note: H = henries = $\Omega*s$
- π = pi
- δ_s = skin depth (m)
- ρ = resistivity ($\Omega*m$)
- ω = radian frequency = $2\pi*f$ (Hz)
- σ = conductivity (mho/m), note: mho [] = siemen [S]

Copper at 10 GHz

$$\delta_s = \sqrt{\frac{2}{(2\pi * 10^{10} \frac{1}{sec}) * (4\pi * 10^{-7} \frac{H}{m}) * (5.8 * 10^7 \frac{mho}{m})}} = \sqrt{\frac{1}{23.2\pi^2 * 10^{10} \frac{H*mho}{m^2*sec} * \frac{\Omega*sec}{H}}} = 6.61 * 10^{-7} m$$

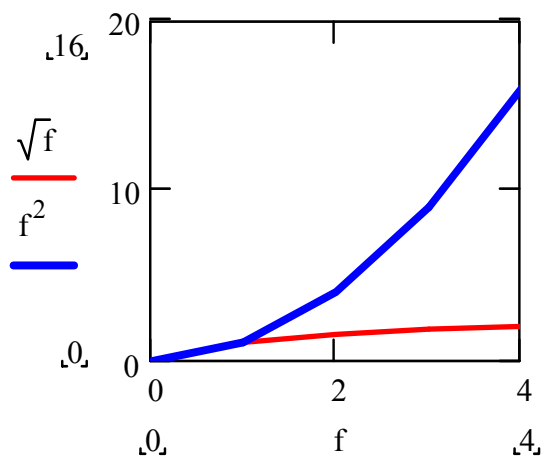
~0.66 μm

15

Material parameters for some SC

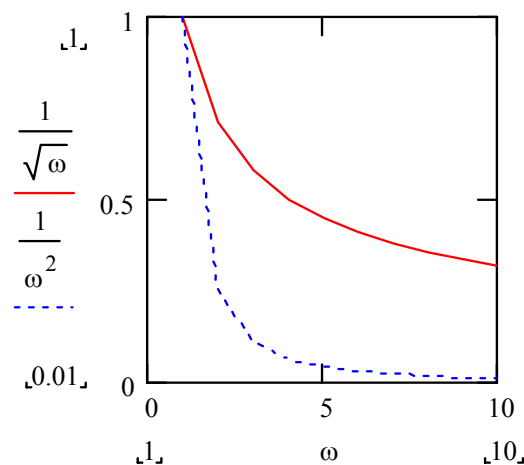
Material	Coherence Length (nm)	London Penetration depth (nm)
Sn	230	34
Al	1600	16
Pb	83	37
Cd	760	110
Nb	38	39

needs only ~40 nm of Nb film!



$$R_s^{SC} \propto f^2$$

$$R_s^{NC} \propto \sqrt{f}$$



$$Q^{NC} \propto \omega^{-1/2}$$

$$Q^{SC} \propto \omega^{-2}$$



Electrical resistivity

- Resistivity ρ is due to scattering
- Scattering rate inversely proportional to the scattering time τ

$$\rho \propto \text{scattering rate} \propto 1/\tau$$

- Matteiesson's rule – scattering rates add

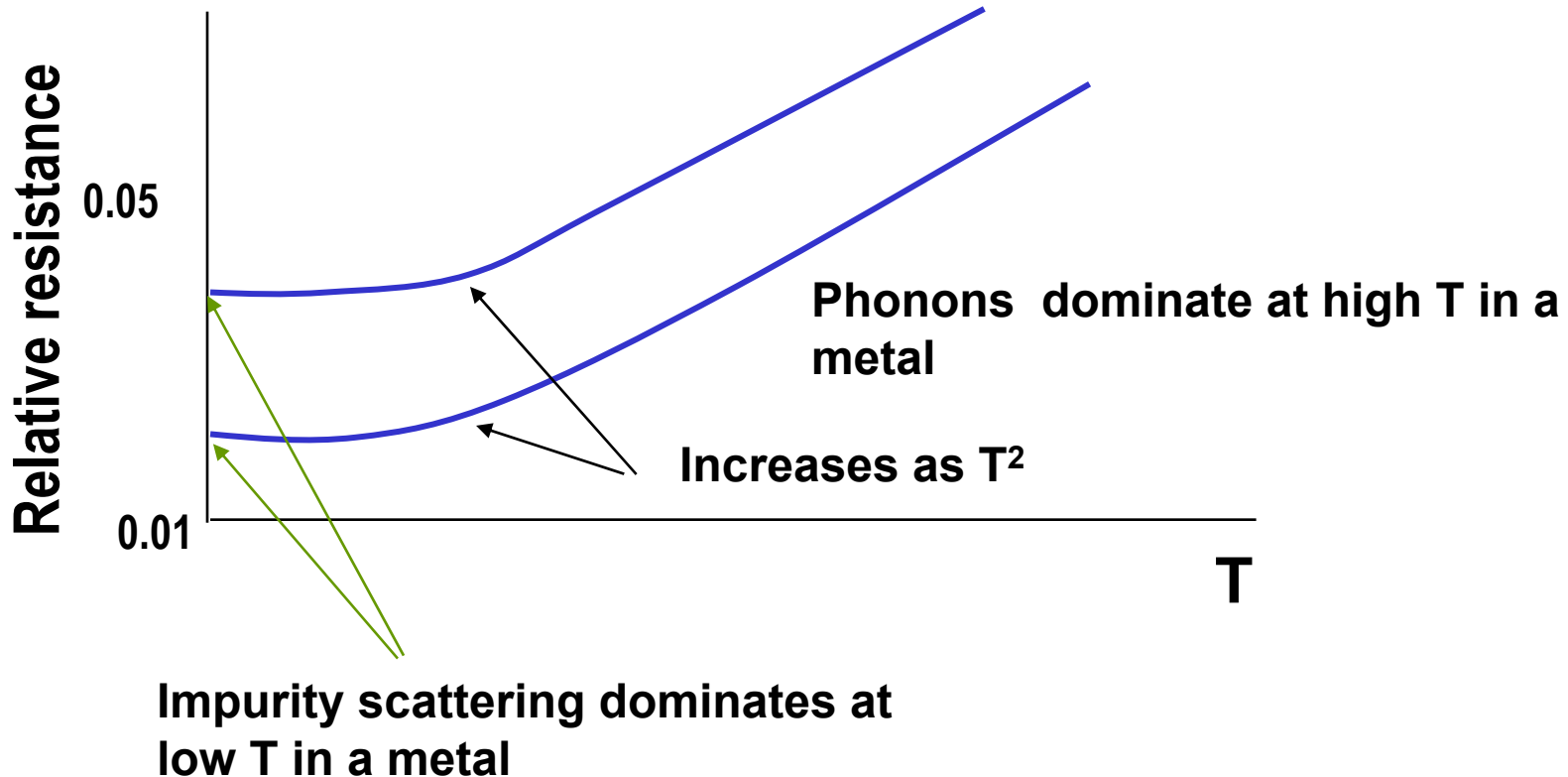
$$\rho = \rho_{\text{vibration}} + \rho_{\text{impurity}} \propto 1/\tau_{\text{vibration}} + 1/\tau_{\text{impurity}}$$

Temperature dependent

Temperature independent

Electrical resistivity

- Consider relative resistivity $R(T)/R(300\text{ K})$
- Typical behavior (for example potassium)



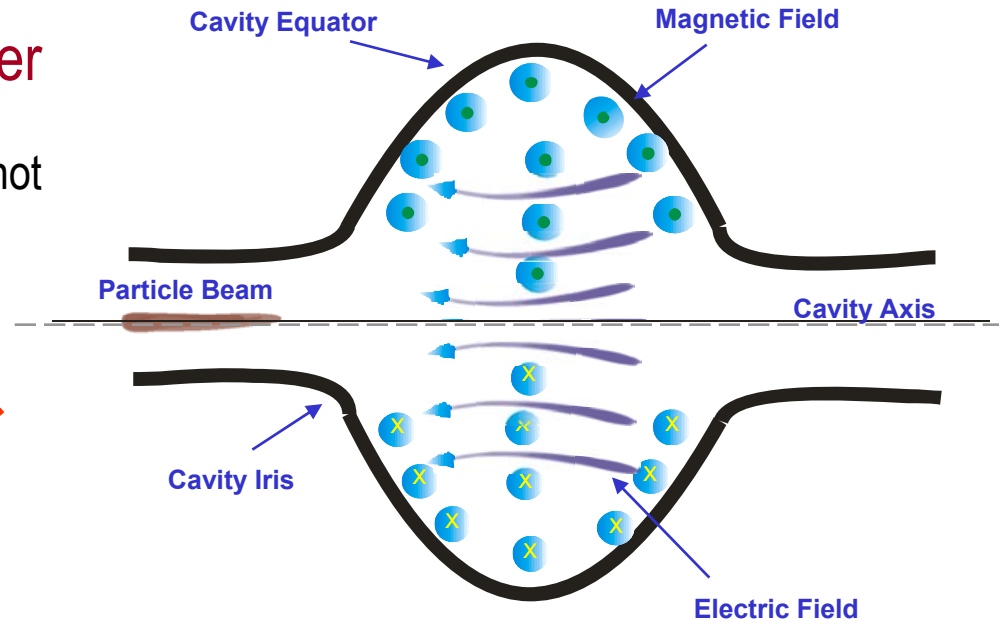
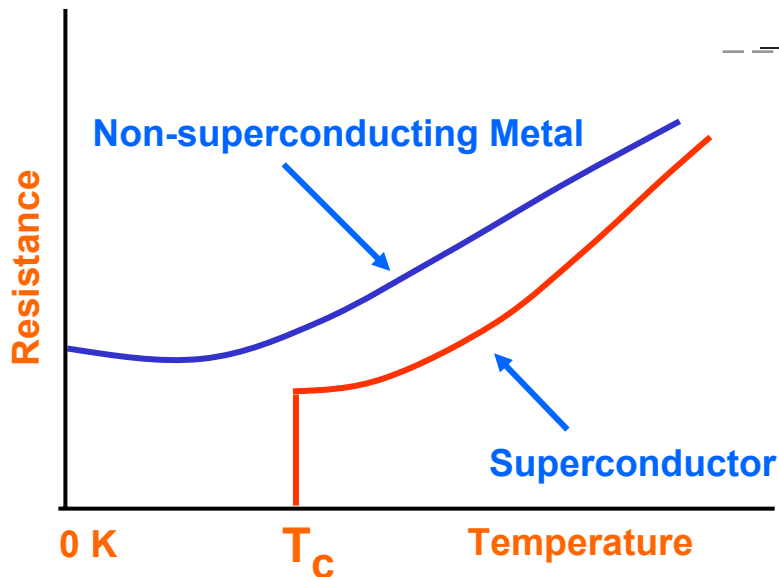
Why use superconducting cavities?

Power dissipation at typical stored energies
(few joules, CW)

● Copper cavity **several MW/meter!**

● Niobium cavity **watts/meter**

Cooling down normal conductors does not help, due to the anomalous skin effect.



Resistivity

$$R_{BCS} \propto f^2 \exp\left(-\frac{\Delta}{KT}\right)$$

At frequencies above 1 GHz the BCS surface resistance at 4.2 K is so high that the gradient in the cavity will be limited by global thermal heating.

The BCS surface resistance can be lowered by reducing the operating temperature.

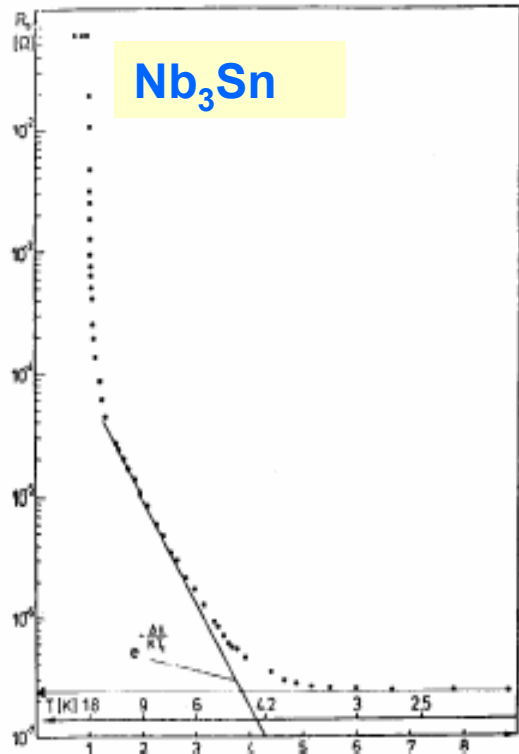
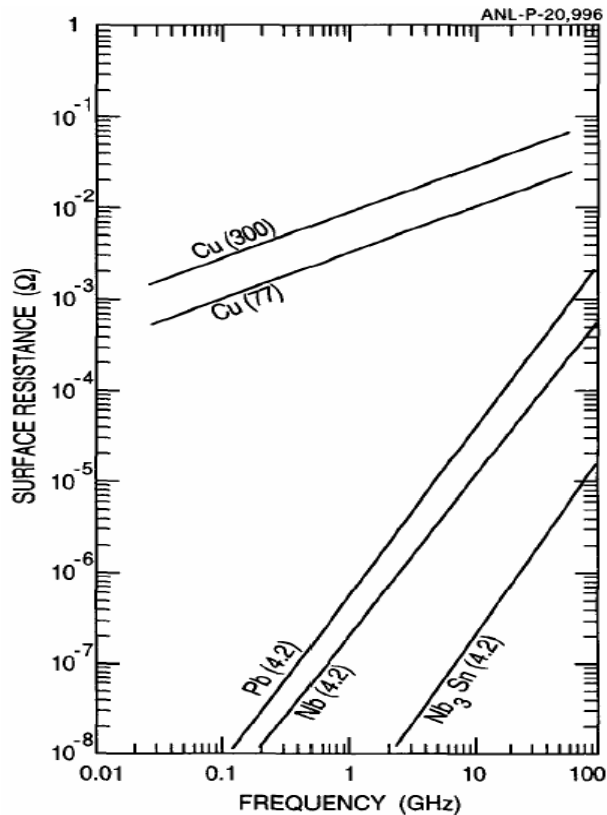
1.8 K is chosen in most cases as the operating temperature for cavities in this frequency range to have some safety margin against crossing the point

$$P_{total} = P_{static} + P_{RF} = P_{static} + P_{BCS} + P_{residual}$$

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Resistivity



$$R_s = R_{BCS} + R_{res}$$

$$R_{BCS} \propto f^2 \exp\left(-\frac{\Delta}{KT}\right)$$

Definition of unloaded Q:

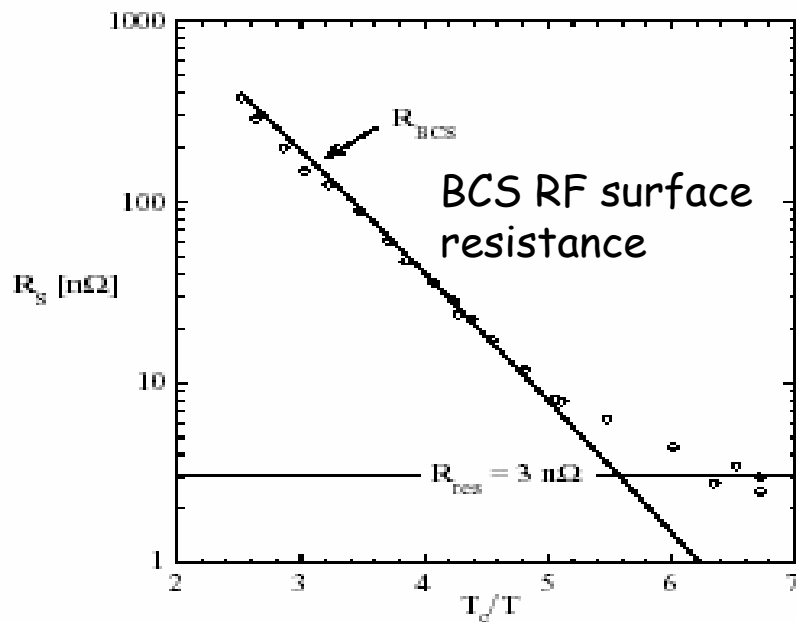
$$Q_o = (wU/P_d) \sim (1/R_s)$$

U= stored energy

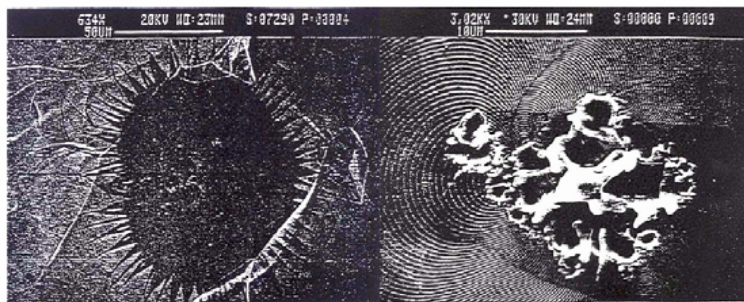
P= power dissipated

R_s=surface resistance

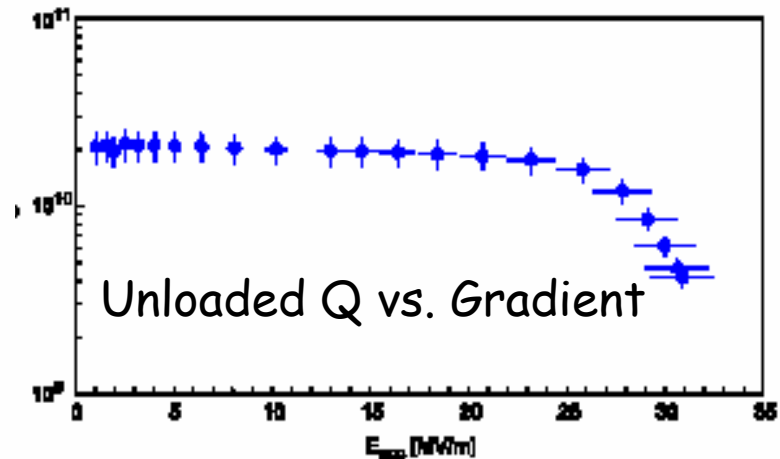
Circumstantial limits



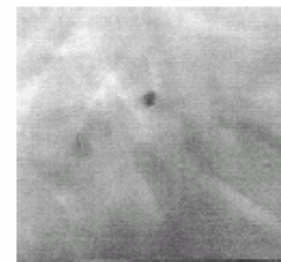
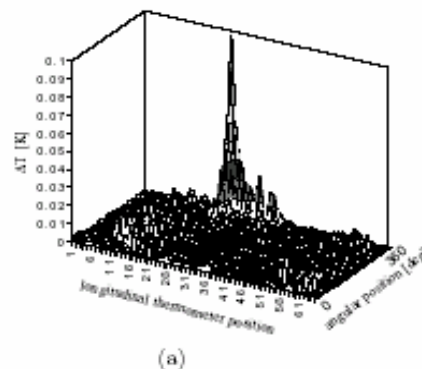
Residual resistance



- Field emission
- Multipacting



"Field" Emission limitations



- Frozen flux
- Solid materials vs. films

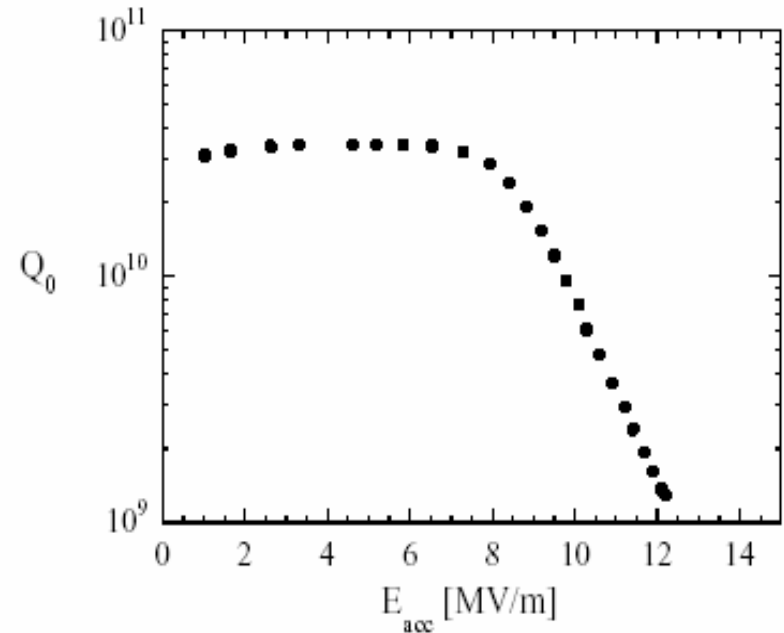
Accelerating Gradient Limitations

● Thermal instability (quench)

- avoiding the normal conducting defect by extreme care in preparing and cleaning the cavity surface;
- increase of the thermal conductivity of the superconducting material.

● Field emission

- Presence of small particles i=on the surface (dusts)
- Cryo absorption of gases
- Acceleration of dark current



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Accelerating Gradient Limitations

● Multipacting

- Resonance multiplication of electrons under the influence of rf field

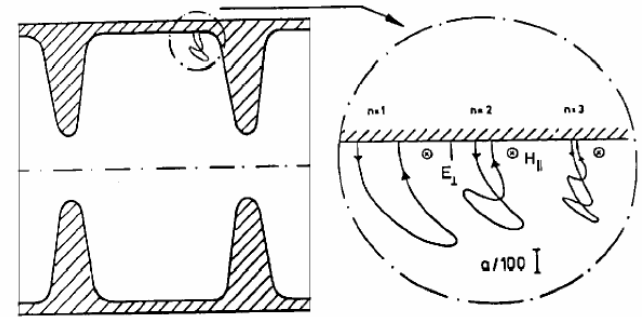
There is a threshold of the field strength in the radiofrequency component. Above this threshold the radiofrequency power can be raised but it has no effect on the stored energy.

The threshold is of sharp nature. The component can be operated below without any sign of unusual behaviour.

The likelihood of multipacting differs for different metals and surface conditions.

Eventually multipacting can be overcome by conditioning. This means that the component behaves normally again after being operated for some time under conditions of multipacting.

Sometimes sparking is observed when operating in the multipacting regime and damage of the component might thereby result.



APS SR (NC)RF

$$P_{total} = P_{beam} + P_{disp} + P_{ol}$$

Ignore this

$$P_{total} = 690kW + 503kW = 1193kW$$

This is the problem

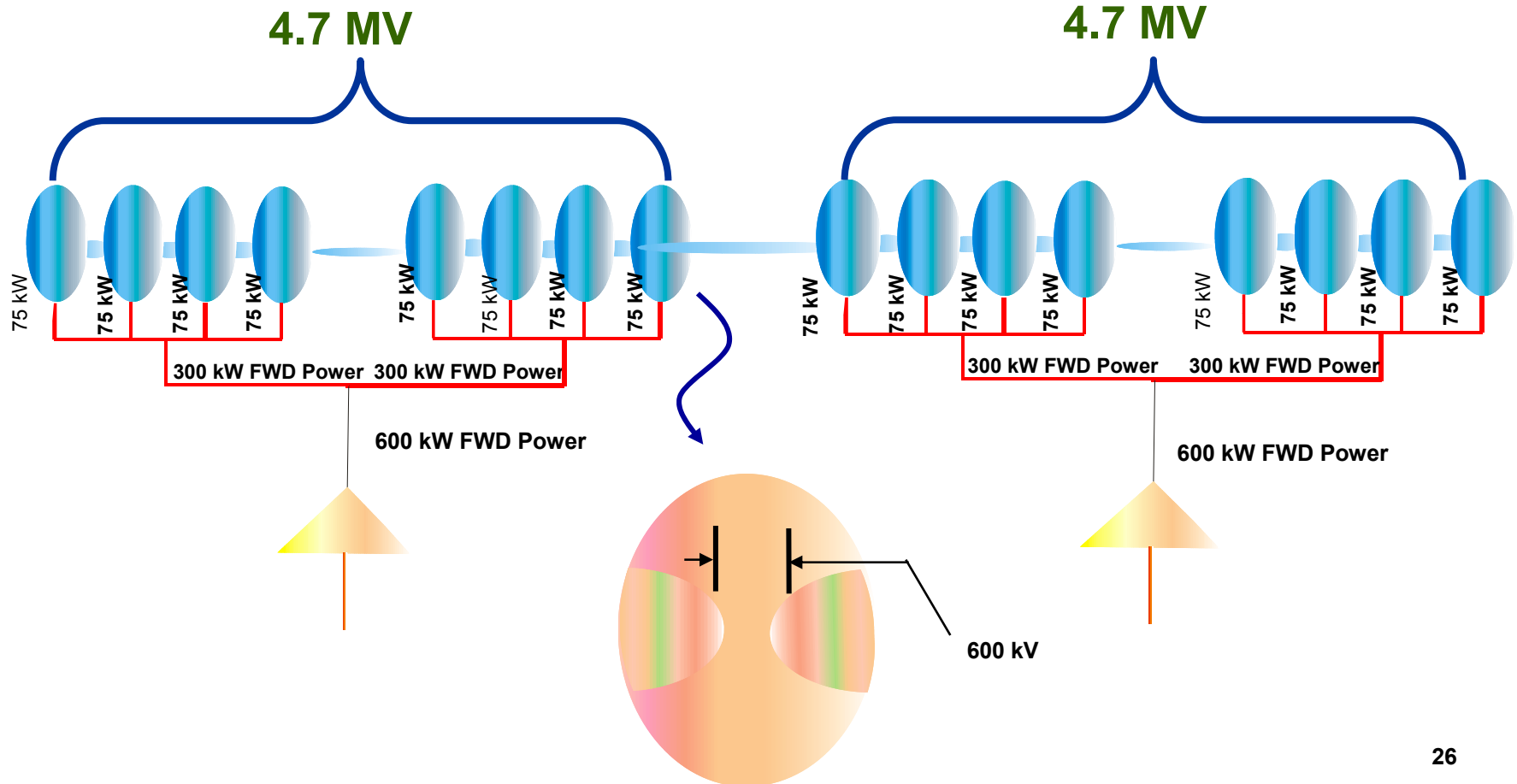
To keep 100 mA, 7 GeV beam circulating in the APS storage ring, we need to provide roughly 1.2 MW CW rf power (@352 MHz):

- Two klystrons, each producing roughly 600 kW rf power
- Each klystron powers eight cavities (2 sectors)

APS SR (NC)RF

Nominal Operating Parameters:

352 MHz, 100 mA @7 GeV, 9.5 MV

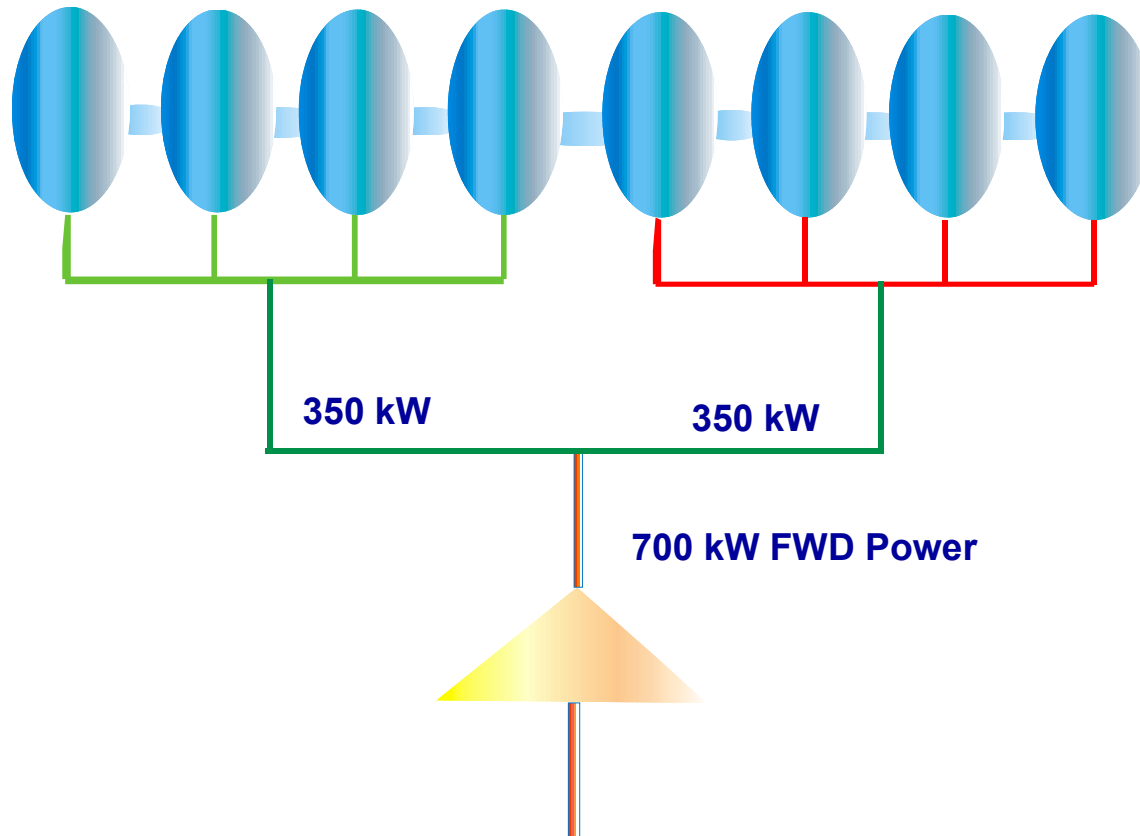


APS SR (SC)RF

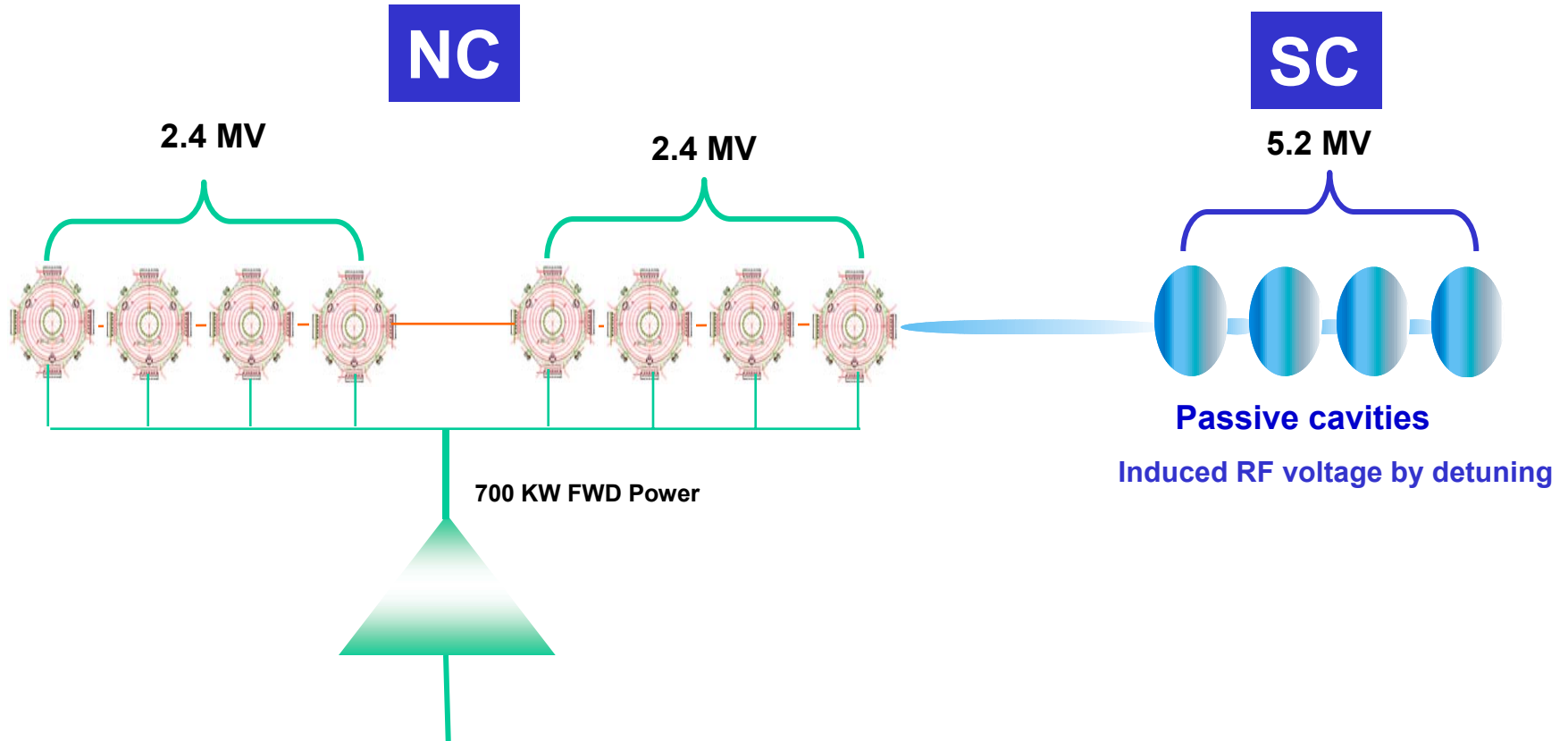
$$P_{total} = P_{beam} + P_{disp} + P_{ol}$$

Ignore this

$$P_{total} = 690kW + 650W = 690.65kW$$



APS SR (Hybrid) RF



Next Lecture

SC RF Cavities March 2nd

- Design criteria
 - frequency
 - power consumption
 - accelerating gradient
- Geometrical considerations and optimization
- Modeling
- Some examples

